Collect. Math. 53, 3 (2002), 303-311
(c) 2002 Universitat de Barcelona

# Mappings with dilatation in Orlicz spaces 

Jana Björn<br>Department of Mathematics, Lund University<br>P. O. Box 118, SE-221 00 Lund, Sweden<br>E-mail: jabjo@maths.lth.se

Received January 25, 2002. Revised April 11, 2002


#### Abstract

We prove openness and discreteness for nonconstant mappings belonging to $W_{\text {loc }}^{1, n}\left(\Omega, \mathbb{R}^{n}\right), n \geq 3$, with dilatation in certain Orlicz spaces which are strictly larger then all $L_{\text {loc }}^{p}(\Omega), p>n-1$. This result contributes to decreasing the gap between known results and a conjecture of Iwaniec and Šverák.


## 1. Introduction

Let $\Omega$ be a connected open set in $\mathbb{R}^{n}, n \geq 2$, and $F: \Omega \rightarrow \mathbb{R}^{n}$ a mapping in the Sobolev space $W_{\text {loc }}^{1,1}\left(\Omega, \mathbb{R}^{n}\right)$. The mapping $F$ is said to have finite dilatation if there exists a function $K$ such that for almost all $x \in \Omega, 1 \leq K(x)<\infty$ and

$$
|D F(x)|^{n} \leq K(x) J_{F}(x),
$$

where $|D F(x)|=\sup \{|D F(x) \xi|:|\xi| \leq 1\}$ is the operator norm of the differential $D F(x)$ of $F$ at $x$ and $J_{F}(x)=\operatorname{det} D F(x)$.

An important theorem due to Reshetnyak [9] states that if $F \in W_{\mathrm{loc}}^{1, n}\left(\Omega, \mathbb{R}^{n}\right)$ is nonconstant and its dilatation $K$ is bounded (i.e. $F$ is quasiregular), then $F$ is continuous, open and discrete, i.e. preimages and images of open sets are open and preimages of single points consist of isolated points. For more about quasiregular mappings, see e.g. Reshetnyak [10] and Rickman [11].

In 1993, Iwaniec and Šverák [4] conjectured that the conclusion of Reshetnyak's theorem is true whenever $F \in W_{\text {loc }}^{1, n}\left(\Omega, \mathbb{R}^{n}\right)$ and $K \in L_{\text {loc }}^{n-1}(\Omega)$, and proved it for $n=2$. Their method does not directly generalize to higher dimensions. An example

Keywords: Capacity, discrete mapping, finite dilatation, Hausdorff measure, open mapping, Orlicz space, Young function.

MSC2000: Primary: 30C65, Secondary: 46E30, 73C50.
by Ball [1] shows that Reshetnyak's theorem fails if we only assume that $K \in L_{\text {loc }}^{p}(\Omega)$ for all $p<n-1$.

By improving Reshetnyak's method, Heinonen and Koskela [2] verified the conjecture of Iwaniec and Šverák for $K \in L_{\mathrm{loc}}^{p}(\Omega), p>n-1$, under the additional assumption that $F$ is quasilight, i.e. that the preimage of each point is compact. The assumption of quasilightness was removed by Villamor and Manfredi [13].

Iwaniec-Koskela-Onninen [5] and Kauhanen-Koskela-Malý [6] weakened the assumption $F \in W_{\text {loc }}^{1, n}\left(\Omega, \mathbb{R}^{n}\right)$ while keeping $K \in L_{\text {loc }}^{p}(\Omega), p>n-1$. The borderline case $p=n-1$ was recently treated by Hencl and Malý [3] by a different method under the assumption of quasilightness.

In this paper, we combine the approach of Kauhanen, Koskela and Maly [6] with a refinement of Villamor and Manfredi's proof and verify the conjecture of Iwaniec and Šverák for a class of mappings with dilatation in Orlicz spaces, namely we prove the following theorem. (See the next section for the definitions of Young functions and Orlicz spaces.)

## Theorem 1

Let $\Psi$ be a doubling Young function such that $\Psi(t) / t^{n-1}$ is nondecreasing, $n \geq 3$, and

$$
\begin{equation*}
\int_{1}^{\infty}\left(\frac{t^{n-1}}{\Psi(t)}\right)^{1 / n(n-2)} \frac{d t}{t}<\infty \tag{1}
\end{equation*}
$$

If $F \in W_{\text {loc }}^{1,1}\left(\Omega, \mathbb{R}^{n}\right)$ is a nonconstant mapping with dilatation belonging to the Orlicz space $L_{\mathrm{loc}}^{\Psi}(\Omega)$, such that

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0+} \varepsilon \int_{\Omega}|D F(x)|^{n-\varepsilon} d x=0 \tag{2}
\end{equation*}
$$

then $F$ is continuous, open and discrete.

Remark 2. For $L^{p}$ spaces (i.e. $\Psi(t)=t^{p}$ ), the condition (1) is satisfied if and only if $p>n-1$. Other Young functions for which the theorem holds are e.g.

$$
\begin{aligned}
& \Psi(t)= \begin{cases}t^{n-1}(\log t)^{n(n-2)+\varepsilon}, & t \geq e \\
t^{n-1}, & 0 \leq t<e\end{cases} \\
& \Psi(t)= \begin{cases}t^{n-1}(\log t)^{n(n-2)}(\log \log t)^{n(n-2)+\varepsilon}, & t \geq e^{e} \\
e^{n(n-2)} t^{n-1}, & 0 \leq t<e^{e}\end{cases}
\end{aligned}
$$

and other repeated logarithms with $\varepsilon>0$. If $\varepsilon=0$ in the above expressions, then the condition (1) just fails.

## 2. Young functions and auxiliary results

Definition 3. A positive continuous convex function $\Psi$ on $(0, \infty)$ satisfying

$$
\begin{equation*}
\lim _{t \rightarrow 0+} \frac{\Psi(t)}{t}=0 \quad \text { and } \quad \lim _{t \rightarrow \infty} \frac{\Psi(t)}{t}=\infty \tag{3}
\end{equation*}
$$

is called a Young function. If $\Psi(2 t) \leq C \Psi(t)$ for some constant $C$ and all $t \in(0, \infty)$, then $\Psi$ is said to be doubling (or satisfying the $\Delta_{2}$-condition).

Definition 4. The Orlicz space $L^{\Psi}(\Omega)$ is the set of all measurable functions with the Luxemburg norm

$$
\|f\|_{L^{\Psi}(\Omega)}=\inf \left\{\lambda>0: \int_{\Omega} \Psi\left(\frac{|f(x)|}{\lambda}\right) d x \leq 1\right\}<\infty
$$

where we interpret $\Psi(0)=0$.
The following generalized Hölder inequality for Orlicz spaces is proved in the paper by O'Neil [7, Theorem 2.3].

## Theorem 5

Let $\Psi_{1}, \Psi_{2}$ and $\Psi_{3}$ be Young functions such that

$$
\Psi_{1}^{-1}(s) \Psi_{2}^{-1}(s) \leq \Psi_{3}^{-1}(s) .
$$

If $f \in L^{\Psi_{1}}(\Omega)$ and $g \in L^{\Psi_{2}}(\Omega)$, then $f g \in L^{\Psi_{3}}(\Omega)$ and

$$
\|f g\|_{L^{\Psi_{3}}(\Omega)} \leq 2\|f\|_{L^{\Psi_{1}}(\Omega)}\|g\|_{L^{\Psi_{2}}(\Omega)} .
$$

The following lemma shows that doubling Young functions can be assumed to be as smooth as needed. Note that for a Young function $\Psi$, the monotonicity of the function $t \mapsto \Psi(t) / t$ is a direct consequence of the convexity of $\Psi$.

## Lemma 6

Let $\Psi \in C^{k}(0, \infty), k \geq 0$, be a positive doubling function satisfying (3), such that the function $t \mapsto \Psi(t) / t$ is nondecreasing. Let

$$
\Psi_{1}(t)=\int_{0}^{t} \frac{\Psi(s)}{s} d s
$$

Then $\Psi_{1}$ is a doubling Young function, $\Psi_{1} \in C^{k+1}(0, \infty)$ and $\Psi(t) / C \leq \Psi_{1}(t) \leq \Psi(t)$ for all $t \in(0, \infty)$. Moreover, $L^{\Psi}(\Omega)=L^{\Psi_{1}}(\Omega)$ and for all $f \in L^{\Psi}(\Omega),\|f\|_{L^{\Psi_{1}}(\Omega)} \leq$ $\|f\|_{L^{\Psi}(\Omega)} \leq 2\|f\|_{L^{\Psi_{1}}(\Omega)}$.

Proof. The fact that $\Psi_{1} \in C^{k+1}(0, \infty)$ follows directly from the definition of $\Psi_{1}$. The convexity of $\Psi_{1}$ is a direct consequence of the fact that $\Psi(t) / t$ is nondecreasing. As for the doubling condition, we have

$$
\Psi_{1}(2 t)=\int_{0}^{t} \frac{\Psi(2 s)}{s} d s \leq C \int_{0}^{t} \frac{\Psi(s)}{s} d s=C \Psi_{1}(t)
$$

To prove (3), let $\varepsilon, k>0$ be arbitrary and find $t_{1}, t_{2}>0$ so that $\Psi(s) / s<\varepsilon$ for all $0<s \leq t_{1}$ and $\Psi(s) / s>k$ for all $s \geq t_{2}$. Then

$$
\begin{array}{ll}
\frac{\Psi_{1}(t)}{t}=\frac{1}{t} \int_{0}^{t} \frac{\Psi(s)}{s} d s \leq \varepsilon & \text { for } 0<t \leq t_{1} \\
\frac{\Psi_{1}(t)}{t} \geq \frac{1}{t} \int_{t / 2}^{t} \frac{\Psi(s)}{s} d s \geq \frac{k}{2} & \text { for } t \geq 2 t_{2}
\end{array}
$$

Next, the fact that $\Psi(t) / t$ is nondecreasing and the doubling property of $\Psi_{1}$ give

$$
\begin{align*}
& \Psi_{1}(t) \geq \int_{t / 2}^{t} \frac{\Psi(s)}{s} d s \geq \frac{t}{2} \frac{\Psi\left(\frac{1}{2} t\right)}{\frac{1}{2} t}=\Psi\left(\frac{1}{2} t\right) \geq \frac{\Psi(t)}{C}  \tag{4}\\
& \Psi_{1}(t)=\int_{0}^{t} \frac{\Psi(s)}{s} d s \leq t \frac{\Psi(t)}{t}=\Psi(t) \tag{5}
\end{align*}
$$

Finally, if $\lambda>\|f\|_{L^{\Psi_{1}}(\Omega)}$, then (4) implies

$$
1 \geq \int_{\Omega} \Psi_{1}\left(\frac{f(x)}{\lambda}\right) d x \geq \int_{\Omega} \Psi\left(\frac{f(x)}{2 \lambda}\right) d x
$$

and hence $2 \lambda \geq\|f\|_{L^{\Psi}(\Omega)}$. Taking infimum over all possible $\lambda$ shows that $\|f\|_{L^{\Psi}(\Omega)} \leq$ $2\|f\|_{L^{\Psi_{1}}(\Omega)}$. Similarly, the inequality $\|f\|_{L^{\Psi_{1}}(\Omega)} \leq\|f\|_{L^{\Psi}(\Omega)}$ follows from (5).

## Lemma 7

Let $\Psi$ be a continuous doubling function on $(0, \infty)$, such that $\Psi(t) / t^{p}$ is nondecreasing. Let $0<\alpha<1-1 / p$ and define $\Psi_{1}$ by $\Psi_{1}^{-1}(s)=s^{\alpha} \Psi^{-1}(s)$. Then $\Psi_{1}$ is a continuous doubling function on $(0, \infty)$ satisfying (3). Moreover, the function $t \mapsto \Psi_{1}(t) / t$ is increasing.

Proof. The continuity is clear. A simple calculation shows that the doubling condition for $\Psi$ is equivalent to $2 \Psi^{-1}(s) \leq \Psi^{-1}(C s)$ for some $C>1$ and all $s \in(0, \infty)$. We then have

$$
2 \Psi_{1}^{-1}(s)=2 s^{\alpha} \Psi^{-1}(s) \leq s^{\alpha} \Psi^{-1}(C s) \leq \Psi_{1}^{-1}(C s)
$$

i.e. $\Psi_{1}$ is doubling. To prove (3), note that

$$
\begin{equation*}
\frac{\Psi_{1}(t)}{t}=\frac{s}{\Psi_{1}^{-1}(s)}=\frac{s^{1-\alpha}}{\Psi^{-1}(s)}=\left(\frac{\Psi(u)}{u^{p}}\right)^{1-\alpha} u^{(1-\alpha) p-1} \tag{6}
\end{equation*}
$$

where $s=\Psi_{1}(t)$ and $u=\Psi^{-1}\left(\Psi_{1}(t)\right) \rightarrow 0+$, as $t \rightarrow 0+$. Hence,

$$
\lim _{t \rightarrow 0+} \frac{\Psi_{1}(t)}{t}=\lim _{u \rightarrow 0+}\left(\frac{\Psi(u)}{u^{p}}\right)^{1-\alpha} u^{(1-\alpha) p-1}=0
$$

and similarly for $t \rightarrow \infty$. Finally, (6) also shows that $\Psi_{1}(t) / t$ is increasing.
We shall also need the following Sobolev type inequality. For the readers convenience, we repeat the short proof.

## Lemma 8

Let $B$ and $B_{0} \subset B \subset \mathbb{R}^{n}$ be balls, $n \geq 3, p>1$ and $1<q<n p /(n-p)$. Then there exists $C>0$ such that for all $u \in W^{1, p}(B)$,

$$
\|u\|_{L^{q}(B)} \leq C\left(\|\nabla u\|_{L^{p}(B)}+\|u\|_{L^{p}\left(B_{0}\right)}\right)
$$

Proof. Assume that this is not true. Then there exist $u_{k} \in W^{1, p}(B)$ so that $\left\|u_{k}\right\|_{L^{q}(B)}=1$ and $\left\|\nabla u_{k}\right\|_{L^{p}(B)}+\left\|u_{k}\right\|_{L^{p}\left(B_{0}\right)} \leq 1 / k$. By the local weak sequentional compactness of $W^{1, p}(B)$, we can find a subsequence, also denoted $u_{k}$, so that $u_{k} \rightarrow u_{0}$ weakly in $W^{1, p}(B)$. The weak lower semicontinuity of the $L^{p}$-norm implies $\left\|\nabla u_{0}\right\|_{L^{p}(B)}=\left\|u_{0}\right\|_{L^{p}\left(B_{0}\right)}=0$ and hence $u_{0}=0$. On the other hand, since the embedding $W^{1, p}(B) \subset L^{q}(B)$ is compact (by e.g. Theorem 2.5.1 in Ziemer [14]), a subsequence of $\left\{u_{k}\right\}_{k=1}^{\infty}$ converges to $u_{0}$ in $L^{q}(B)$ and $\left\|u_{0}\right\|_{L^{q}(B)}=1$. This contradicts $u_{0}=0$.

## 3. Proof of Theorem 1

Definition 9. Let $\Psi$ be a Young function and $K \subset B\left(x_{0}, R\right)=\left\{x \in \mathbb{R}^{n}:\left|x-x_{0}\right|<\right.$ $R\}$ be compact. The $\Psi$-capacity of $K$ with respect to the ball $B\left(x_{0}, 2 R\right)$ is

$$
\operatorname{cap}_{\Psi}\left(K, B\left(x_{0}, 2 R\right)\right)=\inf \|\nabla \phi\|_{L^{\Psi}\left(B\left(x_{0}, 2 R\right)\right)}
$$

where the infimum is taken over all continuous functions $\phi \in W^{1,1}\left(B\left(x_{0}, 2 R\right)\right)$ with compact support in $B\left(x_{0}, 2 R\right)$ and $\phi \geq 1$ on $K$.

The 1-dimensional Hausdorff measure is denoted $H^{1}$. The following theorem is due to Onninen [8].

## Theorem 10

Let $\Psi$ be a doubling Young function satisfying $\Psi(t) \geq C t^{n-1}$ and

$$
\int_{1}^{\infty}\left(\frac{t}{\Psi(t)}\right)^{1 /(n-2)} d t<\infty
$$

If $K \subset B\left(x_{0}, R\right) \subset \mathbb{R}^{n}, n \geq 3$, is compact, then

$$
\begin{equation*}
H^{1}(K) \leq C \operatorname{cap}_{\Psi}\left(K, B\left(x_{0}, 2 R\right)\right) \tag{7}
\end{equation*}
$$

Remark 11. In the definition of $\Psi$-capacity in Onninen [8], only Lipschitz continuous functions are considered. However, the Lipschitz continuity is not used anywhere in the proof. In fact, one only needs that all $x \in K$ are Lebesgue points of $\phi$ in order to obtain the fundamental inequality

$$
1 \leq C \sum_{j=0}^{\infty} \frac{R}{2^{j}} f_{B\left(x, 10 R / 2^{j}\right)}|\nabla \phi(y)| d y
$$

whose integration with respect to the Frostman measure leads to (7).

Proof of Theorem 1. The proof is based on the ideas of Manfredi and Villamor [13]. Note first, that by Theorem 1.3 in Iwaniec-Koskela-Onninen [5] and Theorem 1.5 in Kauhanen-Koskela-Malý [6], $F$ is continuous and sense preserving. The openness and discreteness of $F$ then follows from the Titus-Young theorem (Theorem A in [12]) if we can show that $H^{1}\left(F^{-1}(b)\right)=0$ for every $b \in \mathbb{R}^{n}$ (so that $F^{-1}(b)$ is totally disconnected).

Exhausting $\Omega$ by countably many compacts and covering each of them by finitely many small balls, we can write $\Omega$ as a countable union of balls $B_{j}=B\left(x_{j}, r_{j}\right)$ so that $\overline{B\left(x_{j}, 2 r_{j}\right)} \subset \Omega$ and $F\left(B_{j}\right) \subset B\left(b_{j}, \frac{1}{2} e^{-e}\right), b_{j} \in \mathbb{R}^{n}$. By the $\sigma$-subadditivity of the Hausdorff measure, it suffices to show that $H^{1}\left(B_{j} \cap F^{-1}(b)\right)=0$.

Let $B=B\left(x_{0}, r\right)$ be one of the balls $B_{j}$. Replacing $F$ by $F-b$, we can assume that $b=0 \in F(B)$ and $F(B) \subset B\left(0, e^{-e}\right)$. In [13], Villamor and Manfredi constructed a family of $C^{2}$-smooth radially symmetric $n$-superharmonic functions $\Phi_{a}$ on $B\left(0, e^{-e}\right)$, $0<a<e^{-e}$, with (among others) the following properties,

$$
\begin{array}{rlrl}
\log (1 / a) \leq & \Phi_{a}(y) & \leq \log (1 / a)+\frac{1}{2}+\log 2 & \\
\text { for }|y| \leq a \\
& \Phi_{a}(y)=\log (1 /|y|) & & \text { for } a \leq|y| \leq e^{-e}  \tag{iii}\\
& \Phi_{a}(y) \geq e & & \text { for }|y| \leq e^{-e}
\end{array}
$$

Then, under the assumption $F \in W_{\mathrm{loc}}^{1, n}\left(\Omega, \mathbb{R}^{n}\right)$, they prove that for every nonnegative function $\eta \in C_{0}^{\infty}(\Omega)$,

$$
\begin{equation*}
\int_{\Omega}\left|\nabla\left(\log \Phi_{a} \circ F\right)(x)\right|^{n} \frac{\eta(x)^{n}}{K(x)} d x \leq C \int_{\Omega} K(x)^{n-1}|\nabla \eta(x)|^{n} d x \tag{8}
\end{equation*}
$$

where $C$ is independent of $a$, see (4.1) in [13]. Under the assumptions $F \in W_{\text {loc }}^{1,1}\left(\Omega, \mathbb{R}^{n}\right)$ and (2) considered here, the estimate (8) is proved in Kauhanen-Koskela-Maly [6], formulas (3.1) and (3.2). In the rest of the proof, all constants $C$ will be independent of $a$, they may, however, depend on $F, B$ and other fixed data.

Let $K$ be a compact subset of $F^{-1}(0) \cap B$ and choose a nonnegative function $\eta \in C_{0}^{\infty}(B)$ so that $\eta \geq 1$ on $K$. Following [13], we consider the functions

$$
V_{a}(x)=\frac{\eta(x)\left(\log \Phi_{a} \circ F\right)(x)}{\log \log (1 / a)}
$$

which are continuous by (iii) and satisfy $V_{a} \geq 1$ on $K$ by (i). Moreover, as $F \in$ $W_{\text {loc }}^{1, n-1}(\Omega)$ (by (2)) and $\log \Phi_{a}$ are Lipschitz (using the $C^{2}$-smoothness of $\Phi_{a}$ and (iii)), we have $V_{a} \in W_{0}^{1, n-1}(B)$ and

$$
\nabla V_{a}(x)=\frac{\nabla\left(\log \Phi_{a} \circ F\right)(x) \eta(x)+\left(\log \Phi_{a} \circ F\right)(x) \nabla \eta(x)}{\log \log (1 / a)}
$$

Let $\Psi_{1}(t)=\Psi\left(t^{n}\right)$. It is easily verified that $\Psi_{1}$ is also a doubling Young function and that $K^{1 / n} \in L^{\Psi_{1}}(B)$. Moreover, the function $\Psi_{1}(t) / t^{n(n-1)}$ is nondecreasing. Let $\Psi_{2}$ be defined by $\Psi_{2}^{-1}(s)=s^{1 / n} \Psi_{1}^{-1}(s)$ and let

$$
\Psi_{3}(t)=\int_{0}^{t} \frac{\Psi_{2}(s)}{s} d s
$$

As $n \geq 3$, Lemmas 6 and 7 imply that $\Psi_{3}$ is a doubling Young function. Let us compute the $L^{\Psi_{3}}(B)$-norm of $\nabla V_{a}$. We have

$$
\begin{equation*}
\left\|\nabla V_{a}\right\|_{L^{\Psi_{3}}(B)} \leq \frac{\left\|\left|\nabla\left(\log \Phi_{a} \circ F\right)\right| \eta\right\|_{L^{\Psi_{3}}(B)}+\left\|\left(\log \Phi_{a} \circ F\right)|\nabla \eta|\right\|_{L^{\Psi_{3}}(B)}}{\log \log (1 / a)} \tag{9}
\end{equation*}
$$

As $\Psi_{3}(t) \leq \Psi_{2}(t)$ by Lemma 6 , we have $s^{1 / n} \Psi_{1}^{-1}(s) \leq \Psi_{3}^{-1}(s)$ and an application of the generalized Hölder inequality (Theorem 5) yields

$$
\begin{equation*}
\left\|\left|\nabla\left(\log \Phi_{a} \circ F\right)\right| \eta\right\|_{L^{\Psi_{3}}(B)} \leq 2\left\|\left|\nabla\left(\log \Phi_{a} \circ F\right)\right| K^{-1 / n} \eta\right\|_{L^{n}(B)}\left\|K^{1 / n}\right\|_{L^{\Psi_{1}}(B)} \tag{10}
\end{equation*}
$$

where $\left\|K^{1 / n}\right\|_{L^{\Psi_{1}}(B)}<\infty$ by the assumption. By (8), we have

$$
\begin{equation*}
\left\|\left|\nabla\left(\log \Phi_{a} \circ F\right)\right| K^{-1 / n} \eta\right\|_{L^{n}(B)} \leq C\left(\int_{B} K(x)^{n-1}|\nabla \eta(x)|^{n} d x\right)^{1 / n} \tag{11}
\end{equation*}
$$

and hence the first norm on the right-hand side in (9) is bounded from above by a constant independent of $a$. It remains to estimate the second term on the right-hand side in (9). Another application of the generalized Hölder inequality (Theorem 5) shows that

$$
\begin{align*}
\left\|\left(\log \Phi_{a} \circ F\right) \mid \nabla \eta\right\| \|_{L^{\Psi_{3}}(B)} & \leq C\left\|\log \Phi_{a} \circ F\right\|_{L^{\Psi_{3}}(B)} \\
& \leq 2 C\left\|\log \Phi_{a} \circ F\right\|_{L^{n}(B)}\|1\|_{L^{\Psi_{1}}(B)} \tag{12}
\end{align*}
$$

where $\|1\|_{L^{\Psi_{1}}(B)}<\infty$. The set $B^{\prime}=\{x \in B: F(x) \neq 0\}$ is open and hence there exist $c>0$ and a ball $B_{0} \subset B^{\prime}$ so that $|F(x)|>c$ for all $x \in B_{0}$. In particular, for $a \leq c$ and $x \in B_{0}$, we have $\left|\log \Phi_{a} \circ F(x)\right|<\log \log (1 / c)$ by (ii). As $n \geq 3$, we have $n<n p /(n-p)$ with $p=n-1$ and Lemma 8 implies

$$
\begin{equation*}
\left\|\log \Phi_{a} \circ F\right\|_{L^{n}(B)} \leq C\left(\left\|\nabla\left(\log \Phi_{a} \circ F\right)\right\|_{L^{n-1}(B)}+\left\|\log \Phi_{a} \circ F\right\|_{L^{n-1}\left(B_{0}\right)}\right) \tag{13}
\end{equation*}
$$

The second term on the right-hand side does not exceed $\log \log (1 / c)\left|B_{0}\right|^{1 /(n-1)}$ and the first term is estimated as in Lemma 4 from [13]. More precisely, let $\tilde{\eta} \in C_{0}^{\infty}\left(B\left(x_{0}, 2 r\right)\right)$ be a nonnegative function such that $\tilde{\eta}=1$ on $B$ and $|\nabla \tilde{\eta}| \leq 2 / r$. Then by the Hölder inequality and (8) we have

$$
\begin{align*}
\| \nabla & \left(\log \Phi_{a} \circ F\right) \|_{L^{n-1}(B)} \\
& \leq\left(\int_{B(x, 2 r)}\left|\nabla\left(\log \Phi_{a} \circ F\right)(x)\right|^{n} \frac{\tilde{\eta}(x)^{n}}{K(x)} d x\right)^{1 / n}\|K\|_{L^{n-1}(B(x, 2 r))}^{1 / n} \\
& \leq \frac{C}{r}\|K\|_{L^{n-1}(B(x, 2 r))} \tag{14}
\end{align*}
$$

Putting together (9)-(14) gives

$$
\left\|\nabla V_{a}\right\|_{L^{\Psi_{3}}(B)} \leq \frac{C}{\log \log (1 / a)}
$$

and letting $a \rightarrow 0+$ yields $\operatorname{cap}_{\Psi_{3}}\left(K, B\left(x_{0}, 2 r\right)\right)=0$ for every compact subset $K$ of $F^{-1}(0) \cap B\left(x_{0}, r\right)$. Theorem 10 then shows that $H^{1}(K)=0$, provided that

$$
\int_{1}^{\infty}\left(\frac{t}{\Psi_{3}(t)}\right)^{1 /(n-2)} d t<\infty
$$

As $\Psi_{2}(t) \leq C \Psi_{3}(t)$ by Lemma 6, this follows from the following lemma and finishes the proof of the theorem.

## Lemma 12

Let $\Psi$ be a Young function satisfying (1) such that $\Psi(t) / t^{n-1}$ is nondecreasing. Let $\Psi_{1}$ and $\Psi_{2}$ be as in the proof of Theorem 1. Then

$$
\begin{equation*}
\int_{1}^{\infty}\left(\frac{t}{\Psi_{2}(t)}\right)^{1 /(n-2)} d t<\infty \tag{15}
\end{equation*}
$$

Proof. By Lemma 6, we can assume that $\Psi \in C^{1}(0, \infty)$ and hence also $\Psi_{1}, \Psi_{2} \in$ $C^{1}(0, \infty)$. We have, using integration by parts,

$$
\begin{align*}
\int_{1}^{R}\left(\frac{t}{\Psi_{2}(t)}\right)^{1 /(n-2)} d t= & {\left[\frac{n-2}{n-1}\left(\frac{t^{n-1}}{\Psi_{2}(t)}\right)^{1 /(n-2)}\right]_{1}^{R} } \\
& +\frac{1}{n-1} \int_{1}^{R}\left(\frac{t}{\Psi_{2}(t)}\right)^{(n-1) /(n-2)} \Psi_{2}^{\prime}(t) d t \tag{16}
\end{align*}
$$

Here, the definitions of $\Psi_{1}$ and $\Psi_{2}$ imply

$$
\begin{equation*}
\Psi_{2}^{-1}(s)=\left(s \Psi^{-1}(s)\right)^{1 / n} \tag{17}
\end{equation*}
$$

and hence

$$
\frac{t^{n-1}}{\Psi_{2}(t)}=\frac{\Psi_{2}^{-1}(s)^{n-1}}{s}=\frac{\Psi^{-1}(s)^{1-1 / n}}{s^{1 / n}}=\left(\frac{u^{n-1}}{\Psi(u)}\right)^{1 / n}
$$

where $s=\Psi_{2}(t)$ and $u=\Psi^{-1}\left(\Psi_{2}(t)\right) \rightarrow \infty$, as $t \rightarrow \infty$. As $\Psi(t) / t^{n-1}$ is nondecreasing, we see that the first term on the right-hand side in (16) remains bounded as $R \rightarrow \infty$.

Next, let again $u=\Psi^{-1}\left(\Psi_{2}(t)\right)$, i.e. $\Psi_{2}^{\prime}(t) d t=\Psi^{\prime}(u) d u$. The formula (17) implies $u=t^{n} / \Psi_{2}(t)$ and hence the change of variables $u=\Psi^{-1}\left(\Psi_{2}(t)\right)$ yields that the integral on the right-hand side in (16) is equal to

$$
\int_{\Psi^{-1}\left(\Psi_{2}(1)\right)}^{\Psi^{-1}\left(\Psi_{2}(R)\right)}\left(\frac{u}{\Psi(u)^{n-1}}\right)^{(n-1) / n(n-2)} \Psi^{\prime}(u) d u
$$

Finally, another integration by parts shows that the last integral is equal to

$$
\left[-n(n-2)\left(\frac{u^{n-1}}{\Psi(u)}\right)^{1 / n(n-2)}\right]_{\Psi^{-1}\left(\Psi_{2}(1)\right)}^{\Psi^{-1}\left(\Psi_{2}(R)\right)}+(n-1) \int_{\Psi^{-1}\left(\Psi_{2}(1)\right)}^{\Psi^{-1}\left(\Psi_{2}(R)\right)}\left(\frac{u^{n-1}}{\Psi(u)}\right)^{1 / n(n-2)} \frac{d u}{u}
$$

As before, the first term remains bounded as $R \rightarrow \infty$ and can be disregarded and the second term remains bounded as $R \rightarrow \infty$, by (1).

Acknowledgement. This research has been partly supported by the Swedish Research Council. I also wish to thank Juha Heinonen for introducing me to this problem.

## References

1. J. Ball, Global invertibility of Sobolev functions and the interpenetration of matter, Proc. Roy. Soc. Edinburgh Sect. A 88 (1981), 315-328.
2. J. Heinonen and P. Koskela, Sobolev mappings with integrable dilatations, Arch. Rational Mech. Anal. 125 (1993), 81-97.
3. S. Hencl and J. Malý, Mappings of dinite distortion: Hausdorff measure of zero sets, Preprint 2001.
4. T. Iwaniec and V. Šverák, On mappings with integrable dilatation, Proc. Amer. Math. Soc. 118 (1993), 181-188.
5. T. Iwaniec, P. Koskela, and J. Onninen, Mappings of finite distortion: monotonicity and continuity, Invent. Math. 144 (2001), 507-531.
6. J. Kauhanen, P. Koskela, and J. Malý, Mappings of finite distortion: discreteness and openness, Arch. Ration. Mech. Anal. 160 (2001), 135-151.
7. R. O'Neil, Fractional integration in Orlicz spaces, I, Trans. Amer. Math. Soc. 115 (1965), 300-328.
8. J. Onninen, Orlicz capacities and Hausdorff measures in metric spaces, (in preparation).
9. Yu.G. Reshetnyak, Space mappings with bounded distortion, Sibirsk. Mat. Zh. 8 (1967), 629-658 (in Russian), Siberian Math. J. 3 (1967), 466-486 (English trans1.).
10. Yu.G. Reshetnyak, Spatial mappings with bounded distortion, Nauka, Moscow, 1982 (in Russian), Amer. Math. Soc., Providence, RI, 1989 (English transl.).
11. S. Rickman, Quasiregular mappings, Springer-Verlag, Berlin, 1993.
12. C.J. Titus and G.S. Young, The extension of interiority, with some applications, Trans. Amer. Math. Soc. 103 (1962), 329-340.
13. E. Villamor and J.J. Manfredi, An extension of Reshetnyak's theorem, Indiana Univ. Math. J. 47 (1998), 1131-1145.
14. W.P. Ziemer, Weakly differentiable functions, Sobolev spaces and functions of bounded variation, Springer-Verlag, New York, 1989.
